# Copper-63 Nuclear Quadrupole Resonance Frequencies and Molecular Geometries of Three-co-ordinate Complexes of Copper(I) Halides with $\boldsymbol{N}$-Alkylimidazolidinethione and Thiazolidinethione Ligands $\dagger$ 

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#### Abstract

The crystal structures of seven three-co-ordinated complexes of copper(1) halides with thione ligands, [ $\mathrm{CuXL}_{2}$ ], where $\mathrm{L}=1,3$-thiazolidine-2-thione and $\mathrm{X}=\mathrm{Cl}$ or $\mathrm{Br}, \mathrm{L}=N$-ethylimidazolidine-2-thione and $\mathrm{X}=\mathrm{Br}, \mathrm{L}=N$-isopropylimidazolidine-2-thione and $\mathrm{X}=\mathrm{Cl}$ or Br , and $\mathrm{L}=N$-propylimidazolidine-2thione and $\mathrm{X}=\mathrm{Cl}$ or I have been determined. In all complexes, the copper lies in a trigonal-planar environment and the thiazolidine or imidazolidine rings show no major distortion from planarity. Owing to the presence of hydrogen bonding involving $\mathrm{H}[\mathrm{N}(2)]$, or $\mathrm{H}[\mathrm{N}(02)]$, and the halogen atom all complexes adopt a $W$-shaped conformation. The ${ }^{63} \mathrm{Cu}$ NQR frequencies of these complexes, together with those of two other similar complexes whose structures were already available, depend, for a given halide, on the $\mathrm{S}-\mathrm{Cu}-\mathrm{S}$ bond angle, a decrease in angle producing an increase in resonance frequency: $v_{\mathrm{cl}}=28.7-0.544 \Delta \theta$ and $v_{\mathrm{Br}}=26.9-0.229 \Delta \theta \mathrm{MHz}$, where $\Delta \theta$ is the difference between the $\mathrm{S}-\mathrm{Cu}-\mathrm{S}$ bond angle and $120^{\circ}$. This behaviour has been rationalized in terms of a partial field-gradient model of the resonance frequencies.


We have previously reported studies of the ${ }^{63} \mathrm{Cu} \mathrm{NQR}$ spectra of $\mathrm{Cu}^{1}$ complexes with a variety of charges, co-ordination numbers and ligands. ${ }^{1-9}$ The ${ }^{63} \mathrm{Cu}$ NQR spectra of the three-co-ordinated complexes, [CuXL ${ }_{2}$ ], of $N$-alkylimidazolidine-2thiones ( L ) with cuprous halides ${ }^{1,2}$ led us to believe that there was a possible correlation between the $\mathrm{L}-\mathrm{Cu}-\mathrm{L}$ bond angle and the resonance frequency for the complex of a given halide. To investigate this point further we have determined the crystal structures of seven complexes of this type: $\mathrm{L}=1,3$-thiazolidine-2-thione, $\mathrm{X}=\mathrm{Cl} 1$ or $\mathrm{Br} \mathbf{2} ; \mathrm{L}=N$-ethylimidazolidine-2-thione, $\mathrm{X}=\operatorname{Br} 5 ; \mathrm{L}=N$-isopropylimidazolidine-2-thione, $\mathrm{X}=\mathrm{Cl} 6$ or Br 7 ; and $\mathrm{L}=N$-propylimidazolidine-2-thione, $\mathrm{X}=\mathrm{Cl} 8$ or I 9. The crystal structures of two other complexes [ $\mathrm{CuClL}_{2}$ ] ( $\mathrm{L}=N$-methylimidazolidine-2-thione 3 or $N$-ethylimidazol-idine-2-thione 4) were already available. ${ }^{10,11}$ These results confirm the existence of the correlation and we report here the results of this study.

## Experimental

Preparations.-The preparations of the ligands and the $\mathrm{Cu}^{1}$ complexes $\left[\mathrm{CuXL}_{2}\right]$ have been reported previously. ${ }^{1,2}$ Crystals suitable for X-ray crystallographic studies were obtained by recrystallization from 1,2-dichloroethane (1, 2, 5 and 9), or acetone ( 6 and 7).

X-Ray Crystallography.-Cell parameters and reflection intensities were measured at room temperature on automatic four-circle diffractometers with graphite-monochromated Mo-

[^0]


|  | X | R |  | X | R |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $\mathbf{3}$ | Cl | Me | $\mathbf{7}$ | Br | $\mathrm{Pr}^{\boldsymbol{\beta}}$ |
| $\mathbf{4}$ | Cl | Et | 8 | Cl | $\mathrm{Pr}^{\Omega}$ |
| 5 | Br | Et | 9 | I | $\mathrm{Pr}^{\Omega}$ |
| $\mathbf{6}$ | Cl | Pr |  |  |  |

$K \alpha$ radiation $(\lambda=0.71069 \AA)$. A summary of the crystal data, intensity measurements and structure refinements is given in Table 1, the atomic coordinates are in Table 2. The structures were solved by direct methods (MULTAN $87^{12}$ ) and refined by full-matrix least squares with the XTAL program. ${ }^{13}$ Atomic scattering factors and anomalous dispersion terms are taken from ref. 14. Data were corrected for Lorentz and polarization effects and for absorption. ${ }^{15}$ Selected geometrical parameters are reported in Table 3.
Additional material available from the Cambridge Crystallographic Data Centre comprises H -atom coordinates, thermal parameters and remaining bond lengths and angles.

NQR Spectra.-The NQR spectra were measured on a Decca super-regenerative spectrometer, frequencies being compared to
Table 1 Summary of crystal data, intensity measurement and structure refinement for the $\mathrm{Cu}^{1}$ halide complexes 1, 2, 5-9*

 $\mathrm{CuCl}_{\left(\mathrm{C}_{6} \mathrm{H}_{12} \mathrm{~N}_{2} \mathrm{~S}\right)_{2}}$ 37.5
Orthorhombic $2.098(2)$
$0.689(3)$
$07.568(4)$
90
565(1)
8
1616
0.18 $0.18 \times 0.24 \times 0.25$
1.44
0.58
n 1.294, 1.411
Stoo Stadi-4
2833
1704音 $1 / \sigma^{2}\left(F_{0}\right)$

 | $\mathrm{CuBr}\left(\mathrm{C}_{5} \mathrm{H}_{10} \mathrm{~N}_{2} \mathrm{~S}\right)_{2}$ |
| :--- |
| 403. | Orthorhombic

Pnma
$14.616(4)$
$7.358(3)$
$14.802(2)$
90
$1602.7(3)$
4
816
$0.15 \times 0.25 \times 0.30$ 1.67
0.55 .076
$.752,2.476$
toe Stadi-4
808
54
Calculated
$22{ }^{2}\left(F_{0}\right)$
1.0020
$.51,-0.48$
.90
$.041,0.034$ ${ }^{*} R=\Sigma| | F_{\mathrm{o}}\left|-\left|F_{\mathrm{c}}\right| / \Sigma\right| F_{\mathrm{o}} \mid ; R^{\prime}=\left[\Sigma\left(w\left|F_{\mathrm{o}}\right|-\mid F_{\mathrm{c}}\right)^{2} / \Sigma w\left|F_{\mathrm{o}}\right|^{2}\right]^{\ddagger}, S=\left[\Sigma\left\{\left[\left(F_{\mathrm{o}}-F_{\mathrm{c}}\right) / \sigma\left(F_{\mathrm{o}}\right)\right]^{2}\right\} /\left(N_{\mathrm{r}}-N_{\mathrm{v}}\right)\right]^{\ddagger}$.

Table 2 Atomic coordinates and population parameters with estimated standard deviations in parentheses for the $\mathrm{Cu}^{1}$ halide complexes 1, 2, 5-7 and 9

| Atom | $X / a$ | $Y / b$ | Z/c | Atom | $X / a$ | $Y / b$ | Z/c |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\left[\mathrm{CuCl}\left(\mathrm{C}_{3} \mathrm{H}_{5} \mathrm{NS}_{2}\right)_{2}\right] 1$ |  |  |  |  |  |  |  |
| Cu | $\frac{1}{2}$ | 0.453 31(9) | $\frac{1}{4}$ | S(01) | 0.0077 (1) | 0.4219 (1) | 0.618 93(5) |
| Cl | $\frac{1}{2}$ | 0.093 8(2) | $\frac{1}{4}$ | N(01) | -0.1387(3) | 0.5840 (4) | 0.582 2(1) |
| S(1) | $0.58697(6)$ | 0.618 2(1) | 0.389 85(7) | $\mathrm{N}(02)$ | -0.108 5(4) | 0.422 2(4) | 0.536 4(2) |
| S(2) | 0.678 45(8) | $0.5121(2)$ | 0.607 32(7) | C(01) | -0.082 2(4) | 0.478 2(5) | 0.577 4(2) |
| N | 0.6098 (2) | 0.2315 (5) | 0.4760 (2) | C(02) | -0.218 8(5) | 0.5968 (6) | 0.5431 (2) |
| C(1) | $0.6209(2)$ | 0.437 0(5) | 0.4838 8(2) | C(03) | -0.183 6(5) | $0.4964(6)$ | 0.507 6(2) |
| C(2) | 0.692 8(3) | 0.245 4(6) | 0.653 3(3) | C(04) | -0.141 0(5) | 0.659 4(5) | 0.6263 (2) |
| C(3) | $0.6401(3)$ | $0.1061(6)$ | 0.567 6(3) | C(05) | -0.219 9(8) | $0.6062(7)$ | 0.662 6(3) |
|  |  |  |  | C(06) | -0.163 2(6) | 0.7958 (6) | 0.614 2(3) |
| $\left[\mathrm{CuBr}\left(\mathrm{C}_{3} \mathrm{H}_{5} \mathrm{NS}_{2}\right)_{2}\right]^{2}$ |  |  |  | $\left[\mathrm{CuBr}\left(\mathrm{C}_{6} \mathrm{H}_{12} \mathrm{~N}_{2} \mathrm{~S}\right)_{2}\right] 7$ |  |  |  |
| Cu | $\frac{1}{2}$ | $0.45282(9)$ | $\frac{1}{4}$ | Br | 0.308 78(8) | 0.374 30(4) | 0.315 32(3) |
| Br | $\frac{1}{2}$ | $0.08176(8)$ | $\frac{1}{4}$ | Cu | $0.33081(9)$ | $0.26110(5)$ | $0.41147(4)$ |
| S(1) | 0.58681 (7) | 0.614 8(1) | 0.390 14(6) | S(1) | $0.1567(2)$ | 0.2640 (1) | $0.50735(8)$ |
| S(2) | 0.682 51(8) | 0.521 2(2) | 0.605 22(7) | N(1) | -0.1570(6) | 0.3527 (3) | 0.5373 (2) |
| N | 0.6103 3(2) | 0.2421 1(5) | 0.479 3(2) | N(2) | -0.0561(6) | 0.3951 (4) | 0.4317 (3) |
| C(1) | $0.6217(2)$ | 0.4413 (5) | 0.484 2(2) | C(1) | -0.026 1(7) | 0.339 5(4) | 0.4909 (3) |
| C(2) | $0.6968(3)$ | 0.263 5(7) | 0.655 3(2) | C(2) | -0.2970(8) | 0.4210 (4) | 0.5078 (4) |
| C(3) | 0.6410 (3) | $0.1215(6)$ | 0.571 6(3) | C(3) | -0.226 5(9) | 0.449 6(4) | 0.434 2(4) |
| $\left[\mathrm{CuBr}\left(\mathrm{C}_{5} \mathrm{H}_{10} \mathrm{~N}_{2} \mathrm{~S}\right)_{2}\right] 5$ |  |  |  | C(4) | -0.165 6(8) | 0.307 5(4) | 0.609 4(4) |
|  |  |  |  | C(5) | $-0.350(1)$ | 0.2571 (5) | 0.6123 (4) |
| Cu | 0.305 56(7) | $\frac{3}{4}$ | 0.335 67(7) | C(6) | -0.129(1) | $0.3748(6)$ | 0.6731 (5) |
| Br | $0.14502(8)$ | $\frac{3}{4}$ | 0.344 31(6) | S(01) | 0.5261 (2) | 0.1410 (1) | $0.41728(9)$ |
| S(1) | 0.3750 (2) | $\frac{3}{4}$ | 0.203 6(2) | $\mathrm{N}(01)$ | 0.825 5(6) | 0.0966 (3) | 0.345 0(2) |
| N(1) | $0.3189(5)$ | $\frac{3}{4}$ | 0.0328 (5) | $\mathrm{N}(02)$ | $0.7032(6)$ | 0.2297 (3) | $0.3132(3)$ |
| N(2) | 0.2093 (6) | $\frac{3}{4}$ | 0.128 0(5) | C(01) | 0.689 9(7) | 0.156 4(4) | 0.3565 (3) |
| C(1) | 0.297 5(7) | $\frac{3}{4}$ | $0.1195(6)$ | C(02) | 0.942 2(9) | $0.1305(5)$ | 0.2890 (4) |
| C(2) | 0.2393 (7) | $\frac{3}{4}$ | -0.024 8(6) | C(03) | 0.847 7(8) | 0.2190 (5) | $0.2615(4)$ |
| C(3) | $0.1606(7)$ | $\frac{3}{4}$ | 0.042 5(6) | C(04) | 0.846 6(9) | $0.0062(5)$ | 0.380 3(3) |
| C(4)* | 0.409(1) | 0.826(2) | $0.0009(8)$ | C(05) | $1.052(1)$ | -0.016 7(5) | 0.398 7(4) |
| C(5)* | 0.440(1) | 0.693(3) | -0.068(1) | $\mathrm{C}(06)$ | 0.755(1) | -0.0672(6) | 0.329 6(5) |
| S(01) | 0.3943 (2) | $\frac{3}{4}$ | 0.4561 (1) | $\left[\mathrm{CuI}\left(\mathrm{C}_{6} \mathrm{H}_{12} \mathrm{~N}_{2} \mathrm{~S}\right)_{2}\right] 9$ |  |  |  |
| $\mathrm{N}(01)$ | $0.3587(6)$ | $\frac{3}{4}$ | 0.632 6(5) |  |  |  |  |
| $\mathrm{N}(02)$ | $0.2413(6)$ | $\frac{3}{4}$ | 0.552 4(5) | I | 0.2278 (2) | 0.861 29(6) | 0.340 42(6) |
| C(01) | $0.3301(6)$ | $\frac{3}{4}$ | $0.5513(5)$ | Cu | 0.2167 (2) | 0.7077 (1) | 0.332 51(9) |
| $\mathrm{C}(02)$ | 0.2889 9(8) | $\frac{3}{4}$ | $0.7001(6)$ | S(1) | 0.219 3(5) | 0.640 4(2) | 0.2074 (2) |
| C(03) | 0.2041 (7) | $\frac{3}{4}$ | 0.643 8(6) | N(1) | 0.204(2) | 0.785 6(8) | 0.1265 (6) |
| C(04)* | 0.453(2) | 0.839(4) | 0.656(1) | N(2) | 0.228(2) | 0.684 9(7) | 0.0412 (6) |
| C(05)* | 0.502(2) | 0.672(4) | 0.680(1) | C(1) | 0.218(2) | 0.7063 (9) | 0.1221 (7) |
| $\left[\mathrm{CuCl}\left(\mathrm{C}_{6} \mathrm{H}_{12} \mathrm{~N}_{2} \mathrm{~S}\right)_{2}\right] 6$ |  |  |  | C(2) | 0.209(2) | 0.823 9(8) | 0.043 0(8) |
|  |  |  |  | C(3) | 0.225(2) | 0.753 9(8) | $-0.0177(7)$ |
| Cu | 0.033 50(5) | 0.221 93(6) | $0.59878(2)$ | C(4)* | 0.300(4) | 0.606(2) | $0.013(1)$ |
| Cl | 0.037 3(1) | 0.169 4(1) | 0.519 90(5) | C(41)* | $0.162(6)$ | 0.601(2) | 0.007(3) |
| S(1) | 0.052 2(1) | $0.0837(1)$ | 0.657 61(5) | C(5)* | 0.174(5) | 0.572(2) | -0.063(2) |
| N(1) | $0.1315(4)$ | -0.149 3(4) | 0.660 4(2) | C(51)* | 0.290(6) | 0.580(2) | -0.056(3) |
| N(2) | $0.1127(4)$ | -0.0773(4) | 0.587 4(2) | C(6) | 0.252(2) | 0.492(1) | -0.0901 (9) |
| $\mathrm{C}(1)$ | $0.1009(4)$ | -0.051 5(4) | 0.6340 (2) | S(01) | 0.206 4(6) | $0.6225(2)$ | 0.445 5(2) |
| C(2) | 0.177 5(5) | -0.248 4(5) | $0.6304(2)$ | $\mathrm{N}(01)$ | 0.230(2) | 0.753 7(7) | 0.5478 (6) |
| C(3) | 0.148 6(5) | -0.205 9(6) | $0.5797(2)$ | N(02) | 0.190(2) | 0.642 6(7) | $0.6169(6)$ |
| C(4) | 0.145 0(7) | -0.148 3(6) | 0.712 8(2) | C(01) | 0.211(2) | 0.674 4(8) | 0.5401 (7) |
| C(5) | 0.263(1) | -0.140 6(9) | $0.7262(3)$ | $\mathrm{C}(02)$ | 0.218(2) | 0.781 6(8) | 0.638 6(8) |
| C(6) | $0.0867(8)$ | -0.253(1) | $0.7360(3)$ | C(03) | 0.209(2) | $0.7015(9)$ | 0.687 5(7) |
|  |  |  |  | C(04) | 0.159(2) | 0.553(1) | 0.6369 (9) |
|  |  |  |  | C(05) | 0.332(2) | 0.512(1) | 0.652(1) |
|  |  |  |  | C(06) | 0.288(3) | 0.422(1) | 0.673(1) |

* Atomic site refined with a population parameter of 0.5 .
harmonics from an internal crystal-controlled oscillator. Temperatures were measured with a Hewlett-Packard 2802 digital thermometer and varied between 77 K and room temperature with an Artronix 5301-E temperature controller.


## Results

NQR Spectra.-The ${ }^{63} \mathrm{Cu}$ NQR spectra and their temperature dependence was observed in the range $77-300 \mathrm{~K}$ and the results fitted to the quadratic equation (1). Phase changes occur

$$
\begin{equation*}
v_{T}=v_{0}+A T+B T^{2} \tag{1}
\end{equation*}
$$

for 5 at 238, for 4 at 104 and for 6 at $100 \mathrm{~K} .{ }^{1}$ The presence of two equally intense lines at room temperature for 9 , indicating the presence of two crystallographically inequivalent $\mathrm{Cu}^{1}$ atoms, is confirmed by the crystal structure.

## Discussion

$X$-Ray Crystallography.-Of all the compounds presented here, three complexes show a crystallographic intramolecular symmetry. The isostructural complexes 1 and 2 (Fig. 1) are located on a $C_{2}$ axis with the Cu and X atoms in special positions (Wyckoff sites 4b). In 5, all atoms, except those of
the ethyl substituents at $\mathrm{N}(1)$ and $\mathrm{N}(01)$, are located on a mirror plane (Wyckoff sites 4c); consequently, the imidazolidine rings and the co-ordination plane are coplanar. It should be noted that for this complex a refinement in space group $\mathrm{Pna2}_{1}$, with the molecule in general position, shows that the chirality/ polarity parameter ${ }^{16}$ converges to 0.5 which clearly confirms that this structure is centrosymmetric.
Three structures show disorder of one or both substituents at $\mathrm{N}(1)$ and/or $\mathrm{N}(01)$. As mentioned above, both atomic sites of the ethyl substituent of 5 are located in general positions, out of the mirror plane containing the other atoms of the complex. The refinement of these atomic sites with occupation parameters of 0.5 leads to a coherent geometry of this part of the molecule


Fig. 1 Structure of the $\mathrm{Cu}^{\mathrm{I}}$ halide complex 1 with ellipsoids at $50 \%$ probability


Fig. 2 Structure of the $\mathrm{Cu}^{1}$ halide complex 5 showing the disordered ethyl substituents at $\mathrm{N}(1)$ and $\mathrm{N}(01)$. Ellipsoids are represented with $50 \%$ probability and the disordered atoms with arbitrarily fixed atomic radii of $0.2 \AA$
(Fig. 2). Both complexes with $n$-propyl substituents at $\mathrm{N}(1)$ and $\mathrm{N}(01)$ (8 and 9) exhibit partial or complete disorder of the propyl moieties. In 8, this disorder is significant and also affects the imidazolidine rings of both independent molecules of the asymmetric unit and therefore we only report here the crystal data (Table 2) and the mean observed geometries about the copper atoms (Table 4). In 9, only one of the propyl substituents was disordered and successfully resolved in a crossed disorder by splitting two atomic sites (Fig. 3).

In all complexes, the copper lies in a trigonal-planar environment [maximum deviation of the co-ordination plane $\operatorname{CuXS}(1) S(01)=0.006 \AA$ for 9$]$ and the thiazolidine or imidazolidine rings show no major distortion from planarity. All complexes adopt a $\mathbf{W}$-shaped conformation with the substituents at $\mathrm{N}(1)$ and $\mathrm{N}(01)$ in a anti position relative to the $\mathrm{Cu}-\mathrm{X}$ bond. This particular conformation is due to the presence of hydrogen bonding between $\mathrm{H}[\mathrm{N}(2)]$, or $\mathrm{H}[\mathrm{N}(02)]$, and the halogen atom (Table 3). All dihedral angles between the mean plane passing through the five-membered rings and the co-ordination planes lie in the range $0-14^{\circ}$, except for 6 where this angle is $37.8^{\circ}$ for one of the imidazolidine rings. Nevertheless hydrogen bonding occurs and the $\mathrm{N}(02)-\mathrm{H}[\mathrm{N}(02)]$ bond direction points to the chlorine atom


Fig. 3 Structure of the $\mathrm{Cu}^{1}$ halide complex 9 showing the partially disordered propyl substituent at $\mathrm{N}(1)$. Ellipsoids are represented with $50 \%$ probability and the disordered atoms with arbitrarily fixed atomic radii of $0.2 \AA$

Table 3 Selected geometrical parameters (distances in $\AA$, angles in ${ }^{\circ}$ ) for $\mathrm{Cu}^{1}$ halide complexes 1, 2, 5-7 and 9

| Compound | 1 | 2 | 5 | 6 | 7 | 9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| X | Cl | Br | Br | Cl | Br | I |
| $\mathrm{Cu}-\mathrm{X}$ | 2.282(1) | 2.4116(8) | 2.366 (2) | 2.246(2) | 2.372(1) | 2.531(2) |
| $\mathrm{Cu}-\mathrm{S}(1)$ | 2.2297(9) | 2.2383(8) | 2.206 (3) | 2.205(2) | 2.220(2) | 2.227 (4) |
| $\mathrm{Cu}-\mathrm{S}(01)$ | - | - | 2.210(2) | 2.230(2) | 2.220 (2) | 2.240 (4) |
| $\mathrm{S}(1)-\mathrm{C}(1)$ | 1.686(3) | 1.686(3) | 1.69(1) | 1.691(5) | $1.699(6)$ | 1.71(1) |
| $\mathrm{S}(01)-\mathrm{C}(01)$ | - | - | $1.696(9)$ | 1.691(5) | $1.690(6)$ | 1.69(1) |
| $\mathrm{X}-\mathrm{Cu}-\mathrm{S}(1)$ | 117.98(3) | 118.05(2) | 120.69(8) | 122.85(6) | 123.34(6) | 122.5(1) |
| $\mathrm{X}-\mathrm{Cu}-\mathrm{S}(01)$ | - | - | 123.13(8) | 118.92 (5) | 124.67(6) | 126.0(1) |
| $\mathrm{S}(1)-\mathrm{Cu}-\mathrm{S}(01)$ | 124.03(4) | 123.90(4) | 116.2(1) | 118.23(6) | 111.98(7) | 111.4(1) |
| $\mathrm{Cu}-\mathrm{S}(1)-\mathrm{C}(1)$ | 107.7(1) | 108.8(1) | 109.9(3) | 109.0(2) | 110.8(2) | 110.7(5) |
| $\mathrm{Cu}-\mathrm{S}(01)-\mathrm{C}(01)$ | - | - | 109.9(3) | 105.2(2) | 109.6(2) | 110.9(5) |
| $\mathrm{S}(1)-\mathrm{C}(1)-\mathrm{S}(2)$ | 120.7(2) | 120.2(2) | - | - | - | - |
| $\mathrm{S}(1)-\mathrm{C}(1)-\mathrm{N}(1)$ | - | - | 123.7(8) | 124.1(4) | 123.8(4) | 125(1) |
| $\mathrm{S}(1)-\mathrm{C}(1)-\mathrm{N}(2)$ | 127.7(2) | 128.3(2) | 127.0(7) | 126.1(4) | 125.7(4) | 126.3(9) |
| $\mathrm{S}(01)-\mathrm{C}(01)-\mathrm{N}(01)$ | - | - | 126.9(8) | 124.5(4) | 125.0(4) | 125(1) |
| $\mathrm{S}(01)-\mathrm{C}(01)-\mathrm{N}(02)$ | - | - | 124.6(6) | 125.1(4) | 125.4(4) | 125.4(9) |
| $\mathrm{X}-\mathrm{Cu}-\mathrm{S}(1)-\mathrm{C}(1)$ | -11.6(1) | -11.2(1) | 0.0 | 14.3(2) | 14.2(2) | -3.0(5) |
| $\mathrm{X}-\mathrm{Cu}-\mathrm{S}(01)-\mathrm{C}(01)$ | - | -- | 0.0 | 36.3(2) | 13.6(2) | -1.2(5) |
| $\mathrm{Cu}-\mathrm{S}(1)-\mathrm{C}(1)-\mathrm{N}(2)$ | 7.6(4) | 5.7(2) | 0.0 | -8.1(5) | -4.7(6) | -5(1) |
| $\mathrm{Cu}-\mathrm{S}(01)-\mathrm{C}(01)-\mathrm{N}(02)$ | - | - | 0.0 | -14.1(5) | $1.5(5)$ | -4(1) |
| $\mathrm{X} \cdots \mathrm{H}[\mathrm{N}(2)]^{a}$ | 2.31(3) | 2.45(3) | 2.30 | 2.44 | 2.67 | 2.55 |
| $\mathrm{X} \cdots \mathrm{N}(2)^{\text {a }}$ | 3.167(3) | 3.281(2) | 3.365 | 3.306 | 3.504 | 3.597 |
| Angle between mean planes ${ }^{b}$ |  |  |  |  |  |  |
| c.p. and fmring | 10.9 | 10.6 | 0.0, 0.0 | 13.8, 37.8 | 13.6, 11.9 | 5.8, 4.4 |
| fmring and fmring | 11.4 | 14.2 | 0.0 | 29.6 | 19.9 | 9.8 |

${ }^{a}$ Mean value for 5, 6,7 and 9. ${ }^{b}$ c.p. $=$ Co-ordination plane $[\mathrm{Cu}, \mathrm{X}, \mathrm{S}(1), \mathrm{S}(01)]$; fmring $=$ thiazolidine or imidazolidine ring.
(Fig. 4). The structure of the corresponding bromide complex 7 is shown in Fig. 5.
$N Q R$ Spectra.-Table 4 summarises the NQR frequencies and the bond angles of complexes 1 to 9 . Inspection of these results reveals at once a strong correlation between the $\mathrm{S}-\mathrm{Cu}-\mathrm{S}$ bond angles and the ${ }^{63} \mathrm{Cu} \mathrm{NQR}$ frequency, whereas although the $\mathrm{S}-\mathrm{Cu}$ and the $\mathrm{X}-\mathrm{Cu}$ bond lengths (Table 3) all show significant variations there does not appear to be any obvious correlation with the NQR data. Fig. 6 shows the correlation for the four chlorides $1,3,4$ and 6 and the three bromides. If we express this correlation in terms of the dependence of the frequency, $v$, on the deviation, $\Delta \theta^{\circ}$, of the $\mathrm{S}-\mathrm{Cu}-\mathrm{S}$ angle from $120^{\circ}$ we obtain the expressions (2) and (3). In view of the lack of

$$
\begin{align*}
& v_{\mathrm{Cl}}=28.7-0.544 \Delta \theta  \tag{2}\\
& v_{\mathrm{Br}}=26.9-0.229 \Delta \theta \tag{3}
\end{align*}
$$

precision of the data for the fifth chloride, $\mathbf{8}$, the results for this compound were not included in the correlation but, as can be seen from Fig. 6, the point for this complex lies close to the regression line.


Fig. 4 Structure of the $\mathrm{Cu}^{1}$ halide complex 6 showing the $\mathrm{N}-\mathrm{H}$ bond direction pointing to the chlorine atom. The atomic radii are arbitrarily fixed in order to clarify the diagram


Fig. 5 Structure of the $\mathrm{Cu}^{1}$ halide complex 7 with ellipsoids at $50 \%$ probability

Table 4 The ${ }^{63} \mathrm{Cu}$ NQR frequencies at 77 and 300 K and bond angles [S-Cu-S] of complexes 1-9

| Complex | X | $v_{298} / \mathrm{MHz}$ | $v_{77} / \mathrm{MHz}$ | $\mathrm{S}-\mathrm{Cu}-\mathrm{S} /{ }^{\circ}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\mathbf{1}$ | Cl | 26.479 | 27.400 | 124.0 |
| $\mathbf{4}$ | Cl | 28.802 | $30.296^{a}$ | 118.5 |
|  |  |  | $30.161^{a}$ |  |
| $\mathbf{6}$ | Cl | 30.021 | $32.722^{b}$ | 118.2 |
| $\mathbf{3}$ | Cl | 29.481 | 30.768 | 119.3 |
| $\mathbf{8}$ | Cl | $28.388^{c}$ | $29.671^{c}$ | $120.1^{c}$ |
| $\mathbf{2}$ | Br | 26.321 | 27.059 | 123.9 |
| $\mathbf{5}$ | Br | 27.756 | $29.393^{d}$ | 116.2 |
| $\mathbf{7}$ | Br | 28.730 | 30.343 | 112.0 |
| $\mathbf{9}$ | I | 26.068 | 27.462 | 111.4 |

${ }^{a}$ Phase transition at $104 \mathrm{~K} .{ }^{b}$ Phase transition at $100 \mathrm{~K} .{ }^{c}$ Average for the two different Cu atoms. ${ }^{d}$ Phase transition at 238 K .

In accord with our previous studies this equation indicates that the frequency of a chloride is noticeably higher than that of the corresponding bromide, but it also indicates the new result that the chlorides are much more sensitive to the bond angle than are the bromides.

With only one data point for the iodides we cannot, of course, obtain a correlation, but we can at least attempt to see if the resonance frequency is consistent with the data for the chlorides and bromides. All the systems we have studied so far indicate that for a given type of copper(I) halide complex the frequencies lie in the order $\mathrm{Cl}>\mathrm{Br}>\mathrm{I}$ and that the difference between the chloride and bromide is approximately the same as between the bromide and the iodide. On this assumption the frequency for the iodide for a bond angle of $120^{\circ}$ is 25.2 MHz and the frequency of 9 is then related to the bond angle by equation (4).

$$
\begin{equation*}
v_{\mathrm{I}}=25.2-0.10 \Delta \theta \tag{4}
\end{equation*}
$$

For what it is worth this coefficient of the angular dependence of the iodide is almost exactly in the same ratio to that of the bromide as is the dependence of the bromide to that of the chloride. The conclusion that the frequencies of the iodides are much less sensitive to the $\mathrm{S}-\mathrm{Cu}-\mathrm{S}$ angle than the chlorides or the bromides is confirmed by the fact that the total spread in frequencies ${ }^{2}$ of five such thione iodides is only 2 MHz whereas seven chlorides and seven bromides show a range of values of 5.0 and 3.8 MHz respectively. ${ }^{1}$ For the purposes of the present discussion we therefore take it as established that the resonance frequencies decrease as the $\mathrm{S}-\mathrm{Cu}-\mathrm{S}$ angle is increased above $120^{\circ}$ and that the angular dependence is in the order $\mathrm{Cl}>\mathrm{Br}>\mathrm{I}$. It has also been established for a wide variety of systems that for a given type of complex the frequencies of the halides vary in the same order. ${ }^{1-9}$

The high polarisability of the $\mathrm{Cu}^{\mathrm{I}}$ cation together with the participation of both 3 d and 4 p electrons in the bonding between the central atom and its ligands makes almost impossible even the simple Townes-Dailey ${ }^{17}$ analysis of ${ }^{63} \mathrm{Cu}$ NQR frequencies. When this difficulty is compounded with the fact that for a nucleus of spin $\frac{3}{2}$ such as ${ }^{63} \mathrm{Cu}$ the singleresonance frequency is given by equation (5) and that neither

$$
\begin{equation*}
v_{Q}=\frac{e^{2} Q q_{z z}}{2 h} \sqrt{1+\frac{\eta^{2}}{3}} \tag{5}
\end{equation*}
$$

the coupling constant, $e^{2} Q q_{z z}$, nor the asymmetry parameter, $\eta$, is known separately, it becomes clear that we must turn to other methods. An alternative technique for accounting for the resonance frequencies of polyco-ordinated nuclei is the partial field-gradient method. ${ }^{18}$ In this method the effect of a given


Fig. 6 The ${ }^{63} \mathrm{Cu} \mathrm{NQR}$ frequencies (in MHz at 298 K ) and $\mathrm{S}-\mathrm{Cu}-\mathrm{S}$ bond angles of three-co-ordinated complexes of copper( I ) chlorides $(\mathrm{O})$, bromides $(\square)$ and iodides $(\nabla)$ with thione ligands


Scheme 1
ligand on the field gradient of the central atom is represented by a partial cylindrically symmetrical field-gradient tensor symmetrical about the bond between the central atom and the ligand. The total field-gradient tensor is then given by the sum of these partial tensors. Although we have already demonstrated that for $\mathrm{Cu}^{1}$ complexes it is not possible to characterize a given ligand by a partial field gradient that can be used for a variety of co-ordination numbers of the central atom and a variety of charges of the complex, ${ }^{2}$ we have used it with success to account for the difference in resonance frequency between two crystalline modifications of a neutral four-co-ordinated $\mathrm{Cu}^{1}$ complex that had slightly different geometries in the two modifications. ${ }^{19}$ It therefore seemed possible that the method could also be useful in this context.

The idealized structure of these complexes is shown in Scheme 1. If the partial field gradients of A and B are represented by $a$ and $b$ respectively, and the A-Cu-A angle by $2 \theta$ then $q_{\mathrm{a}}$ and $q_{\mathrm{b}}$, the components of the field-gradient tensor in the molecular plane perpendicular to and along the $\mathrm{Cu}-\mathrm{B}$ bond, and $q_{\mathrm{c}}$, the component perpendicular to the molecular plane are given by equations (6)-(8). For three identical ligands and $\theta=$

$$
\begin{gather*}
q_{\mathrm{a}}=2 a\left(3 \sin ^{2} \theta-1\right)-b  \tag{6}\\
q_{\mathrm{b}}=2 a\left(3 \cos ^{2} \theta-1\right)+2 b  \tag{7}\\
q_{\mathrm{c}}=-(2 a+b) \tag{8}
\end{gather*}
$$

$60^{\circ}$ the asymmetry parameter is zero and the $q_{z z}$ direction is $q_{\mathrm{c}}$, perpendicular to the molecular plane. When the ligands are no longer identical then, provided that the partial field gradients of the ligands A and B are not too different and that $2 \theta$ deviates little from $120^{\circ}$ the $q_{z z}$ axis will continue to lie along the direction $q_{\mathrm{c}}$. This is the case for all six complexes with this structure whose field-gradient tensors have been completely determined by Zeeman measurements. For other combinations of ligands the $q_{z z}$ axis may lie along $q_{\mathrm{a}}, q_{\mathrm{b}}$ or $q_{\mathrm{c}}$ and as the parameters are varied it may abruptly switch from one direction to another, as indeed will $q_{x x}$ and $q_{y y}$. However, for a spin $\frac{3}{2}$ nucleus such as ${ }^{63} \mathrm{Cu}$ it can be shown that there is an exact relationship between the resonance frequency and the three components of the field-gradient tensor that does not depend either on a knowledge of the asymmetry parameter or of the orientation of the directions of these components with respect to the molecular framework. Thus, rewriting equation (5) explicitly in terms of the three field-gradient components we obtain equations (9)-(11); however, $q_{z z}=-\left(q_{x x}+q_{y y}\right)$ so

$$
\begin{gather*}
v=\frac{e^{2} Q q_{z z}}{2 \sqrt{3}} \sqrt{3+\frac{\left(q_{x x}-q_{y y}\right)^{2}}{q_{z z}^{2}}}  \tag{9}\\
v=\frac{e^{2} Q}{2 \sqrt{3}} \sqrt{3 q_{z z}^{2}+q_{x x}^{2}+q_{y y}^{2}-2 q_{x x} q_{y y}}  \tag{10}\\
\left.v=\frac{e^{2} Q}{2 \sqrt{3}} \sqrt{q_{z z}^{2}+q_{x x}^{2}+q_{y y}^{2}-2\left(q_{x x} q_{y y}-q_{z z}^{2}\right.}\right) \tag{11}
\end{gather*}
$$

that $q_{z z}^{2}=-q_{z z}\left(q_{x x}+q_{y y}\right)$, and therefore equation (12) can be

$$
\begin{equation*}
v=\frac{e^{2} Q}{2 \sqrt{3}} \sqrt{q_{z z}^{2}+q_{x x}^{2}+q_{y y}^{2}-2\left(q_{x x} q_{y y}+q_{y y} q_{z z}+q_{z z} q_{x x}\right)} \tag{12}
\end{equation*}
$$

derived. Since each component of the field-gradient tensor occurs in this formula in exactly the same way-either as a square or as a cross product with both of the other components-we may substitute $q_{\mathrm{a}}, q_{\mathrm{b}}$ and $q_{\mathrm{c}}$ for $q_{x x}, q_{y y}$ and $q_{z z}$ without having to worry about the correspondence between the pairs of tensor components. The resonance frequency is therefore a smooth function of $q_{\mathrm{a}}, q_{\mathrm{b}}$ and $q_{\mathrm{c}}$ and thus of $a, b$ and $\theta$, equations (13) and (14). In the present context we are

$$
\begin{equation*}
v=e^{2} Q \sqrt{q_{\mathrm{a}}^{2}+q_{\mathrm{b}}^{2}+q_{\mathrm{c}}^{2}-2\left(q_{\mathrm{a}} q_{\mathrm{b}}+q_{\mathrm{b}} q_{\mathrm{c}}+q_{\mathrm{c}} q_{\mathrm{a}}\right)} \tag{13}
\end{equation*}
$$

$$
\begin{align*}
& v= \\
& \quad e^{2} Q \sqrt{12 a^{2} \cos ^{2} \theta-\left(12 a^{2}-6 a b\right) \cos ^{2} \theta+4 a^{2}-2 a b+b^{2}} \tag{14}
\end{align*}
$$

interested in the derivative of this expression with respect to $\theta$. This may be straightforwardly obtained as in equation (15).
$\frac{\partial \nu}{\partial \theta}=$
$\frac{e^{2} Q\left[-48 a^{2} \cos ^{3} \theta \sin \theta+\left(24 a^{2}-12 a b\right) \cos \theta \sin \theta\right]}{2 \sqrt{12 a^{2} \cos ^{4} \theta-\left(12 a^{2}-6 a b\right) \cos ^{2} \theta+4 a^{2}-2 a b+b^{2}}}$

When the ligands are identical then for $\theta=60^{\circ} \partial v / \partial \theta=0$, while it can be shown that for this same angle if $a>b$ the derivative is positive and if $a<b$ the derivative is negative. We have already shown that for these complexes the partial field gradients are in the order $\mathrm{Cl}>\mathrm{Br}>\mathrm{I} \approx>\mathrm{C}=\mathrm{S} .{ }^{1,2,20} \mathrm{An}$ approximate value for $e^{2} Q q_{\mathrm{a}}$ appropriate to the thione ligand is 18 MHz while for $\mathrm{Cl}, \mathrm{Br}$ and I the values of $e^{2} Q q_{\mathrm{b}}$ are 26,22 and 18 MHz respectively. These values imply an angular dependence for thione complexes of $-0.206,-0.112$ and 0.0 MHz per degree for $\mathrm{Cl}, \mathrm{Br}$ and I respectively, in reasonable agreement with the experimental values.

## Conclusion

The dependence of the ${ }^{63} \mathrm{Cu} \mathrm{NQR}$ frequency of three-coordinated complexes of copper(I) halides and $N$-alkyl-imidazolidine-2-thione ligands on the $\mathrm{S}-\mathrm{Cu}-\mathrm{S}$ angle has been demonstrated experimentally and a satisfactory explanation of this effect has been found in the partial field-gradient model.

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